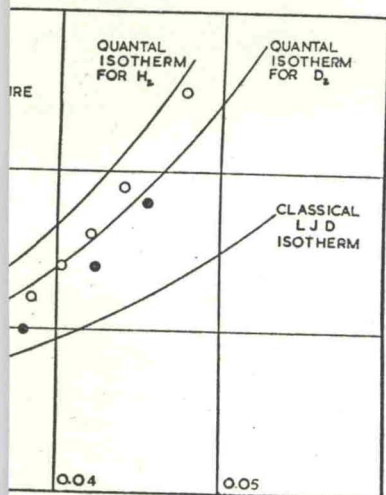


are given in tables 1 and 2 in the form of pressure-density curves, as shown schematically in fig. 2 and 3.



(MOLES .CM³)
 isotherms at 78.9° K. The open circles are for H₂, the filled circles are for D₂.

DISCUSSION

Compared with the classical LJD and quantum isotherms at 78.9° K. The theoretical curves for the potential energy between a pair of molecules are given by eqn. (1) with $\epsilon^* = 50.75 \times 10^{-16}$ ergs, and the virial coefficients of H₂ and D₂. The difference between H₂ and D₂ molecules is just noticeable at gases 7 when the temperature is high enough. We have interpolated in the table of Curtiss³ to find the classical LJD pressures of eqn. (2) to obtain the quantal isotherms at the experimental points lie closer to the quantum isotherms, particularly at the higher densities. This fact alone, however, is not sufficient to explain the quantum inflation of the pressure. It may be inaccurate at these reduced pressures. A significant fact, therefore, is that the difference in density is roughly that predicted by the classical theory. This fact alone, however, is not sufficient to explain the quantum inflation of the pressure. It may be inaccurate at these reduced pressures. A significant fact, therefore, is that the difference in density is roughly that predicted by the classical theory.

The quantum inflation of the pressure can be met by plotting the experimental data together with the results for other similar gases. It is then possible to compare the experimental data for light gases and the predictions of the classical theory. For this, we introduce the quantities:

$$p^* = pv_0/N\epsilon^* = pD^3/\epsilon^*,$$

$$v^* = v/v_0 = v/ND^3,$$

$$T^* = kT/\epsilon^*,$$

where v_0 and D are the van der Waals' constants. It is also convenient to express the results in terms of the dimensionless parameters

where $y^* = h^2 D(m\epsilon^*)^{-1/2}$ which de Boer^{11,12} showed occurs in the general expression for the reduced equation of state in quantum statistics. Eqn. (2) then becomes

$$\Delta p^* = T^* \left[v^{*-1} + \frac{3}{2} \frac{d \log y^*}{d v^*} \right] [9.0710 y^{*1/2} T^{*1/2} v^{*1/2} A^{*-1} - 1]^{-1}, \quad (3)$$

where y^* is a function of v^* , the relationship between them being

$$(1 + 12y^* + 25.2y^{*2} + 12y^{*3} + y^{*4})(1 + y^*)^{-1}(1 - y^*)^{-6} - 2v^{*2} = 0. \quad (4)$$

It is in conformity with the general considerations of de Boer^{11,12} our reduced equation of state is of the type

$$p^* = f(T^*, v^*, A^*),$$

where f is a universal function.

For the general plot of p^* against v^* in fig. 4, we have selected the reduced temperature $T^* = 2.14$ corresponding to our experimental temperature $T = 78.9^\circ \text{K}$

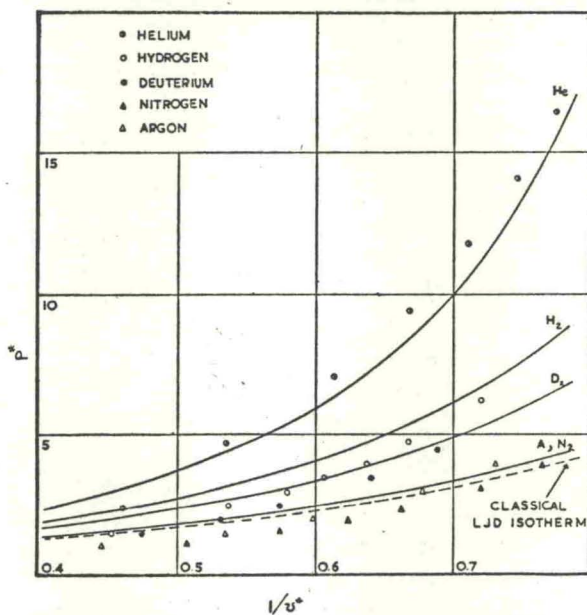


Fig. 4.—Reduced isotherms for $T^* = 2.14$. The full curves are the quantal isotherms.

TABLE 3.—MOLECULAR CONSTANTS

	ϵ^*/k (°K)	D (cm $\times 10^{-8}$)	A^*
He	10.2	2.56	2.64
H ₂	37.0	2.92	1.73
D ₂	37.0	2.92	1.22
N ₂	96.6	3.72	0.225
A	120.3	3.41	0.187

for H₂ and D₂. The intermolecular force parameters and the values of A^* have been taken from the tables of de Boer.^{11,12} They are listed in table 3. The experimental data for He, N₂ and A were given by short extrapolations on the temperature scale from the measurements of Buchmann¹³ on He; Bartlett, Hetherington, Kvalnes and Tremearne¹⁴ on N₂; Michels, Wijker and Wijker¹⁵ on A.